

Invited

Photo-and Electroluminescence Investigation of Rare Earth Ions in III-V Semiconductors

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Sharply structured photoluminescence bands were observed in the III-V semiconductors InP, GaP, and GaAs doped with the rare earth ($4f^n$)-ions Yb, Er, Nd and Pr. These emission bands in the near infrared spectral region are due to internal $4f - 4f$ transitions of trivalent $4f^n$ -ions. Zeeman measurements, optically detected magnetic resonance, electron spin resonance and photoluminescence excitation measurements were performed to determine the structure of the $4f^n$ -centers and to separate different rare earth complexes. The characteristic intra $4f$ -shell transitions could also be excited electrically in $4f$ -doped pn-diodes and laser diodes giving rise to incoherent and coherent radiation around $1.0 \mu\text{m}$ and $1.54 \mu\text{m}$.

1. Introduction

The investigation of the luminescence properties of rare earth ($4f^n$)-ions in solids is an important branch in solid state physics since several decades. This research has concentrated dominantly on ionic hosts such as oxides and fluorides /1/ and to a much lesser extent on II-VI semiconductors /2/.

The most remarkable feature of these ions is the spectral sharpness of their $4f$ -emissions even at room temperature, which results from the shielding of the inner lying $4f$ -shell by the outer filled $5s^2$ and $5p^6$ orbitals. Due to the small interaction of the $4f$ -orbitals with the crystalline environment rare earth ions exhibit characteristic intra $4f$ -shell transitions in the visible and in the near infrared spectral region. The high solubility of $4f$ -ions in insulating hosts is a reason why they were studied extensively as the gain media in externally pumped laser crystals such as $\text{Nd}^{3+}:\text{YAG}$ /1/ and $\text{Er}^{3+}:\text{YAlO}_3$ /3/.

Since a few years the investigations of the luminescence of $4f$ -ions in solids concentrated also on the technologically important III-V semiconductors and silicon, also with

the aim to achieve novel types of optoelectronic devices. This article reviews briefly the spectroscopic investigations of the trivalent $4f^n$ -ions Yb^{3+} ($4f^{13}$), Nd^{3+} ($4f^3$), Pr^{3+} ($4f^2$) and Er^{3+} ($4f^{11}$) in the III-V semiconductors InP, GaP, and GaAs. The feasibility of producing light emitting diodes and laser diodes by different kinds of epitaxial growths processes will also be reviewed.

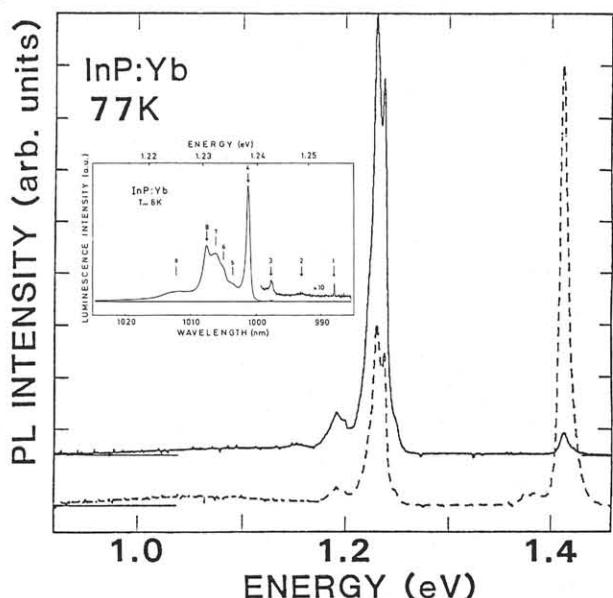


Fig. 1: Luminescence spectrum of InP:Yb grown by MOCVD /15/. The insert shows the high resolution spectrum of Yb^{3+} /17/.

2. Results and Discussion

Ytterbium (Yb) in InP, GaP, and GaAs

The first case of a rare earth ion related emission was observed for Yb-doped InP grown by molten solution /4-6/ and for Yb-doped GaP grown by the same technique /7-9/. This luminescence band was also observed for Yb-implanted InP, GaP, and GaAs samples after proper annealing steps /10-12/. The possibility to grow Yb-doped InP epitaxial layers by liquid phase epitaxy (LPE) and metal organic chemical vapour deposition (MOCVD) was demonstrated by Haydl et al. /13/ and Uwai et al. /14, 15/, respectively; Fig. 1. The observed transition around $1.0 \mu\text{m}$ in these Yb-doped III-V semiconductors is due to the internal $4f - 4f$ transition $2F_{5/2} \rightarrow 2F_{7/2}$ of Yb^{3+} ($4f^{13}$). The fine structure in the low temperature photoluminescence spectra results from the crystal field splitting of the spin-orbit levels; see insert of Fig. 1. The excited states of the luminescent level in InP:Yb were determined by photoluminescence excitation spectroscopy /16/. Zeeman measurements /17/ and optically detected magnetic resonance /18/ revealed that the Yb^{3+} spectrum in InP arises entirely from only one type of cubic Yb^{3+} -center, which resides presumably on a substitutional cation site. Besides this dominant cubic (Td) Yb^{3+} -center an Yb^{3+} -center with trigonal symmetry (line No. 1, see insert of Fig. 1) was identified /11/. Electron spin resonance on InP:Yb was observed and revealed a g-factor for the Γ_6 ground state with $g = 3.291 \pm 0.001$ /19-21/, which is somewhat larger than the one determined by Zeeman measurements /17/.

Neodymium (Nd) in GaP and GaAs

Three sets of sharp emission lines were observed for Nd-implanted GaP due to the intra-center transitions $4F_{3/2} \rightarrow 4I_{9/2}$ (900 nm), $4F_{3/2} \rightarrow 4I_{11/2}$ (1100 nm) and $4F_{3/2} \rightarrow 4I_{13/2}$ (1400 nm) of Nd^{3+} ($4f^3$) /22/. Similar emission spectra were also recorded in Nd-implanted GaAs /22/. Isothermal and isochronal annealing

studies and photoluminescence excitation measurements /23/ revealed that the multiplicity of the observed emission lines arises from different Nd^{3+} -centers with non cubic symmetry /22, 23/. The nature of the defects which destroys the cubic symmetry is not known as yet.

Praseodymium (Pr) in GaP

A rich structured Pr-related emission spectrum could be detected in Pr-doped GaP grown by molten solution. This intra 4f-shell transition of Pr^{3+} ($4f^2$) could only be observed after heat treatment of the sample /8/. If Pr was implanted in GaP and co-implanted with Li different emission spectra were recorded /24/ presumably due to an association of Pr with Li.

Gadolinium (Gd) in InP

No luminescence data exist for Gd^{3+} ($4f^7$) in III-V semiconductors, but electron spin resonance due to Gd^{3+} could be observed in InP /25/ grown by molten solution. The fine structure lines with $g = 1.985 \pm 0.005$ result from a Gd^{3+} -center in the $8S_{7/2}$ ground state exhibiting cubic symmetry. Furthermore, a Gd^{3+} -center with lower than cubic symmetry was detected.

Erbium (Er) in InP, GaP, and GaAs

A luminescence band around $1.54 \mu\text{m}$ due to the internal transition $4I_{13/2} \rightarrow 4I_{15/2}$ of Er^{3+} ($4f^{11}$) was observed in Er-implanted InP, GaP,

MBE GaAs:Er

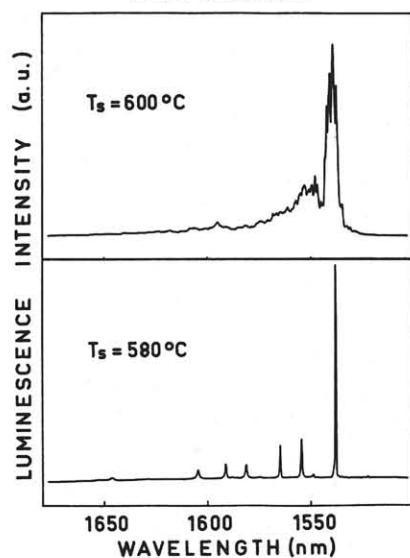


Fig. 2: Effect of substrate temperature T_s on the Er^{3+} -emission at $1.54 \mu\text{m}$. Excitation: 514.5 nm, $T=6 \text{ K}$ /31/.

and GaAs /26-28/. Different annealing conditions result in the creation of several kinds of Er^{3+} -centers in GaAs /28/ presumably due to Er^{3+} -complexes. In Bridgman grown GaAs:Er a luminescence spectrum was recorded with only a few lines which are different from those observed in Er-implanted samples /12/. Er-doping during epitaxial growth like LPE /21,29/ and MBE /30/ results also in quite different emission spectra of Er^{3+} around $1.54 \mu\text{m}$.

For MBE-grown GaAs:Er layers the structure of the luminescence spectrum depends on the substrate temperature /31/; see Fig. 2. The linewidth of the main emission line at $1.542 \mu\text{m}$ (lower part of Fig. 2) is very small and amounts to 0.5 cm^{-1} . Higher resolution spectra taken with a Fourier transform spectrometer (FTIR) reveals a linewidth of 0.04 cm^{-1} /32/ at low temperature. At room temperature the halfwidth is of the order of 10 cm^{-1} . The analysis of the temperature dependent photoluminescence measurements and photoluminescence excitation measurements /31/ on MBE GaAs:Er which exhibits a luminescence spectrum as shown in the lower part of Fig. 2 indicates that this spectrum arises entirely from an Er^{3+} -complex with non cubic symmetry. Zeeman measurements /32/ have shown that the Er^{3+} -center has C_{1h} symmetry. On the other hand Zeeman measurements on LPE-GaAs:Er reveals that this Er^{3+} -center has cubic symmetry /29/.

Electroluminescence

Yb-doped InP layers grown by molten solution /33/ and LPE /13, 34/ were used to prepare light emitting devices. It was shown that the internal 4f - 4f transitions $2\text{F}_{5/2} \rightarrow 2\text{F}_{7/2}$ of Yb^{3+} at $1.0 \mu\text{m}$ could be excited electrically in these pn-diodes. The variation of the electroluminescence intensity of Yb^{3+} with applied voltage indicates that the space charge recombination current might be responsible for the excitation of the Yb^{3+} -emission /13/. The external quantum efficiency for the Yb^{3+} -emission

is of the order of 10^{-3} at 77 K . Light emitting diodes prepared from MBE-GaAs:Er /35/ and MBE-Si:Er /36/ show emission spectra around $1.54 \mu\text{m}$ due to the internal 4f - 4f transition $4\text{I}_{13/2} \rightarrow 4\text{I}_{15/2}$ of Er^{3+} ($4f^{11}$). All the above mentioned 4f-related emissions in III-V semiconductors and silicon were observed at wavelengths which are longer than those of the band edge emission of the host semiconductors. A new type of a 4f-doped laser structure was published by Tsang et al. /37/ and van der Ziel et al. /38/. In this hetero-epitaxial ridge overgrown laser structure the band edge emission wavelength of the quaternary III-V semiconductor was somewhat longer than the Er^{3+} emission wavelengths. These InGaAsP:Er injection lasers where the Er-ions are imbedded in the active layer, exhibit single longitudinal mode operation at $1.5322 \mu\text{m}$ which can already be observed at 300 K /37/, see Fig. 3. This lasing line shifted at a slow rate of $1 \text{ \AA}/\text{^{\circ}C}$ with heat sink temperature /37/ which has to be compared to the drastic shift of the emission wavelengths ($6.5 \text{ \AA}/\text{^{\circ}C}$) of the non Er-doped quaternary semiconductor lasers.

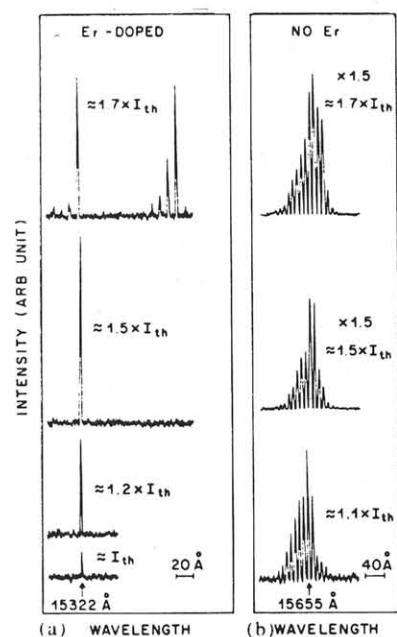


Fig. 3: Lasing spectra of an Er-doped HRO laser and a non Er-doped HRO laser at different current injection levels /37/.

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